



Dynamics and Real-Time Simulation (DARTS) Laboratory

### **Spatial Operator Algebra (SOA)**

10. Closed-Chain Dynamics (Cut-Joint Method)

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https://dartslab.jpl.nasa.gov/



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# SOA Foundations Track Topics (serial-chain rigid body systems)



- 1. Spatial (6D) notation spatial velocities, forces, inertias; spatial cross-product, rigid body transformations & properties; parallel axis theorem
- 2. Single rigid body dynamics equations of motion about arbitrary frame using spatial notation
- **3. Serial-chain kinematics** minimal coordinate formulation, hinges, velocity recursions, Jacobians; first spatial operators; O(N) scatter and gather recursions
- **4. Serial-chain dynamics** equations of motion using spatial operators; Newton–Euler mass matrix factorization; O(N) inverse dynamics; composite rigid body inertia; forward Lyapunov equation; mass matrix decomposition; mass matrix computation; alternative inverse dynamics
- **5.** Articulated body inertia Concept and definition; Riccati equation; alternative force decompositions
- **6. Mass matrix factorization and inversion** spatial operator identities; Innovations factorization of the mass matrix; Inversion of the mass matrix
- Recursive forward dynamics O(N) recursive forward dynamics algorithm; including gravity and external forces; inter-body forces identity



## **SOA Generalization Track Topics**



- 8. Graph theory based structure BWA matrices, connection to multibody systems
- 9. Tree topology systems generalization to tree topology rigid body systems, SKO/SPO operators, gather/scatter algorithms
- **10.Closed-chain dynamics (cut-joint)** holonomic and non-holonomic constraints, cut-joint method, operational space inertia, projected dynamics
- **11.Closed-chain dynamics (constraint embedding)** constraint embedding for graph transformation, minimal coordinate closed-chain dynamics
- **12.Flexible body dynamics** Extension to flexible bodies, modal representations, recursive flexible body dynamics





## Recap



## **Previous Session Recap**



- Built upon BWA concepts to define SKO and SPO operators for multibody systems
- Defined the general class of SKO-models for multibody systems
- Showed the virtually all the analysis and algorithms developed for serial-chain, rigi-body systems carries over to SKO models with only minor generalizations
- This opens the door for applying the operator methods and algorithms to any multibody system with an SKO model
  - As we will see, this is a very broad class of multibody systems





## **Closed-Chain Dynamics**



## **Closed-chain multibody systems**

- Our development so far has focused on multibody systems that have SKO models
  - Serial-chain systems
  - Branched tree systems
- The underlying bodies topology has that been of a tree
- Closed-chain system topologies have loops, and hence do not have tree structure or an obvious SKO model







#### **Example Closed-Chain Dynamics Mechanisms**

Time step: 0.001s

0.00s





Time step: 0.001s

0.00s



## **Suspension system example**





Multiple closed-loops



## **Closed Chain Modeling Options**





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- Bilateral constraints: Defined by an equality relationship
  - eg. mechanism loops
  - always active when present
  - Hinges are actual a dual way of describing a bilateral constraint
- Unilateral constraints: Defined by inequality constraint
  - eg. contact/collision dynamics
  - can be active/inactive based on the current state



## **Holonomic Bilateral Constraints**



Smooth constraint on the coordinates

$$\mathfrak{d}(\theta,t)=\mathbf{0}$$

Reduces dofs from  $\mathcal{N}$  to  $(\mathcal{N} - n_c)$  dimension.

Differentiating:

$$\begin{split} \mathbf{\hat{\vartheta}}(\theta,t) &= G_{c}(\theta,t)\mathbf{\dot{\theta}} - \mathfrak{U}(t) = \mathbf{0} \\ \\ G_{c}(\theta,t) &\triangleq \nabla_{\theta}\mathfrak{d}(\theta,t) \in \mathcal{R}^{n_{c} \times \mathcal{N}} \end{split} \qquad \begin{aligned} \mathfrak{U}(t) &\triangleq -\frac{\partial \mathfrak{d}(\theta,t)}{\partial t} \in \mathcal{R}^{n_{c}} \end{aligned}$$





# The constraints are expressed directly at the velocity level:

$$\boldsymbol{\dot{\vartheta}}(\boldsymbol{\theta},t) = G_{c}(\boldsymbol{\theta},t)\boldsymbol{\dot{\theta}} - \mathfrak{U}(t) = \boldsymbol{0}$$





Using Lagrange multipliers, DAE form of the equations of motion is:

$$\begin{split} \mathcal{M}(\theta)\boldsymbol{\ddot{\theta}} + \mathcal{C}(\theta,\boldsymbol{\dot{\theta}}) - G_{c}^{*}(\theta,t)\lambda &= \mathcal{T} \\ G_{c}(\theta,t)\boldsymbol{\dot{\theta}} &= \mathfrak{U}(t) \end{split}$$

The dimension of the Lagrange multipliers and the row dimension of  $G_c$  increases with increase in number of cut-joints.

The Lagrange multipliers are the inter-body constraint forces for the cut-joints.



**Rearranged matrix form** 



Using Lagrange multipliers, DAE form

$$\begin{split} \mathcal{M}(\theta)\boldsymbol{\ddot{\theta}} + \mathcal{C}(\theta,\boldsymbol{\dot{\theta}}) - G_{c}^{*}(\theta,t)\lambda &= \mathfrak{T} \\ G_{c}(\theta,t)\boldsymbol{\dot{\theta}} &= \mathfrak{U}(t) \end{split}$$

have

$$\begin{pmatrix} \mathcal{M} & \mathbf{G}_{c}^{*} \\ \mathbf{G}_{c} & \mathbf{0} \end{pmatrix} \begin{bmatrix} \mathbf{\ddot{\theta}} \\ -\lambda \end{bmatrix} = \begin{bmatrix} \mathcal{T} - \mathcal{C} \\ \mathbf{\acute{U}} \end{bmatrix}$$

 $\acute{\mathfrak{U}} \stackrel{\Delta}{=} \overset{\Delta}{\mathfrak{U}}(t) - \overset{\bullet}{\mathbf{G}}_{c} \overset{\bullet}{\boldsymbol{\theta}} \in \mathcal{R}^{n_{c}} \qquad \overset{\bullet}{\mathfrak{d}}(\boldsymbol{\theta}, t) \stackrel{11.3,11.6}{=} G_{c} \overset{\bullet}{\boldsymbol{\theta}} - \acute{\mathfrak{U}}$ 





## **Solution Approaches**



## 1. Projection solution method

- Switch to minimal coordinates form
- Pick  $(N n_c)$  of the coordinates as independent variables
- Numerically project the equations of motion down to these independent variables
- Solve these equations of motion and lift up the solution to get all coordinate accels
- Expensive process, and has issues with picking indep coords









The main idea is to numerically project the dynamics down to a minimal set of coordinates and hence eliminate the explicit constraints

$$G_{c}X_{c} = \mathbf{0}$$
  $\mathbf{\dot{\theta}} = \mathbf{\dot{\theta}}_{p} + X_{c}\mathbf{\dot{\theta}}_{r}$  reduced minimal coordinates

$$\mathfrak{T}-\mathfrak{C}=\mathfrak{M}\boldsymbol{\ddot{\theta}}-\mathsf{G}_{c}^{*}\boldsymbol{\lambda}\overset{11.34}{=}\mathfrak{M}(\underline{\boldsymbol{\ddot{\theta}}}_{p}+X_{c}\boldsymbol{\ddot{\theta}}_{r})-\mathsf{G}_{c}^{*}\boldsymbol{\lambda}$$

$$\mathcal{M}_{r} \stackrel{\triangle}{=} X_{c}^{*} \mathcal{M} X_{c} \in \mathcal{R}^{\mathcal{N}-n_{c} \times \mathcal{N}-n_{c}} \begin{bmatrix} \text{projected} \\ \text{mass} \\ \text{matrix} \end{bmatrix}$$

 $\mathcal{M}_{\mathbf{r}}\mathbf{\bar{\theta}}_{\mathbf{r}} = \mathbf{X}_{\mathbf{c}}^{*}(\mathcal{T} - \mathcal{C} - \mathcal{M}\underline{\mathbf{\ddot{\theta}}}_{\mathbf{p}})$ 

projected equations of motion





- The projected methods is a minimal coordinates, and hence ODE approach
- However, the mass matrix is obtained by a numerical projection approach – which destroys all structure, and we are left with an expensive to compute mass matrix, with opaque structure
- The lack of structure means that SKO models are not applicable and the recursive techniques cannot be used



## 2. Direct solution method

• Set up and solve this equation numerically

$$\begin{pmatrix} \mathcal{M} & \mathbf{G}_{c}^{*} \\ \mathbf{G}_{c} & \mathbf{0} \end{pmatrix} \begin{bmatrix} \mathbf{\ddot{\theta}} \\ -\lambda \end{bmatrix} = \begin{bmatrix} \mathcal{T} - \mathcal{C} \\ \mathbf{\acute{\mu}} \end{bmatrix}$$

- Usually used by absolute coordinate approaches which use maximal cuts so have individual bodies
  - The "tree" system mass matrix is constant & sparse and consists of independent bodies







## 3. Augmented solution method

• Use <u>minimal</u> number of joint cuts so have a spanning tree + cut-joint constraints

$$\begin{pmatrix} \mathcal{M} & \mathbf{G}_{c}^{*} \\ \mathbf{G}_{c} & \mathbf{0} \end{pmatrix} \begin{bmatrix} \mathbf{\ddot{\theta}} \\ -\lambda \end{bmatrix} = \begin{bmatrix} \mathcal{T} - \mathcal{C} \\ \mathbf{\acute{U}} \end{bmatrix}$$

• The tree system is a minimal coordinate multibody system with a configuration dependent mass matrix







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### 4. Constraint embedding solution approach

- Structure based minimal coordinate approach
- Uses graph transformation and variable geometry bodies
- Will cover later ...







## **Augmented Solution Method**



### **Augmented solution approach**



### Have

$$\begin{pmatrix} \mathfrak{M} & \mathbf{G}_{c}^{*} \\ \mathbf{G}_{c} & \mathbf{0} \end{pmatrix} \begin{bmatrix} \mathbf{\ddot{\theta}} \\ -\lambda \end{bmatrix} = \begin{bmatrix} \mathfrak{T} - \mathfrak{C} \\ \mathfrak{U} \end{bmatrix}$$

## **Inverse expression**

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix}^{-1} = \begin{pmatrix} \mathbf{I} & -A^{-1}B \\ \mathbf{0} & \mathbf{I} \end{pmatrix} \begin{pmatrix} A^{-1} & \mathbf{0} \\ \mathbf{0} & F_2^{-1} \end{pmatrix} \begin{pmatrix} \mathbf{I} & \mathbf{0} \\ -CA^{-1} & \mathbf{I} \end{pmatrix}$$
$$= \begin{pmatrix} A^{-1} + A^{-1}BF_2^{-1}CA^{-1} & -A^{-1}BF_2^{-1} \\ -F_2^{-1}CA^{-1} & F_2^{-1} \end{pmatrix}$$

$$\mathsf{F}_2 = (\mathsf{D} - \mathsf{C}\mathsf{A}^{-1}\mathsf{B})$$



## **Rearranged solution equations**







## **Simplified dynamics equations**









- 1. Use O(N) method to solve  $\ddot{\theta}_{f} \stackrel{\Delta}{=} \mathcal{M}^{-1}(\mathcal{T} \mathcal{C})$ "free" generalized accels
- 2. Compute  $G_c \mathcal{M}^{-1} G_c^*$  and use to solve for the Lagrange multipliers in  $\lambda = - [G_c \mathcal{M}^{-1} G_c^*]^{-1} \mathfrak{F}_f$
- 3. Use O(N) method to solve  $\ddot{\boldsymbol{\theta}}_{\delta} \triangleq \mathcal{M}^{-1}G_{c}^{*} \lambda$  for correction generalized accels
- 4. Combine the free and correction generalized accels in  $\ddot{\theta} = \ddot{\theta}_{f} + \ddot{\theta}_{\delta}$

Non-minimal coordinates, DAE approach. Still can use SKO methods for the spanning tree SKO model



### **Mass matrix – but singular**



$$\begin{split} \mathbf{\ddot{\theta}} &= \mathbf{\mathcal{Y}}_{C} \left[ \mathbf{\mathcal{T}} - \mathbf{\mathcal{C}} \right] + \mathbf{\mathcal{M}}^{-1} \mathbf{G}_{c}^{*} \left[ \mathbf{G}_{c} \mathbf{\mathcal{M}}^{-1} \mathbf{G}_{c}^{*} \right]^{-1} \mathbf{\acute{\mu}} \\ \\ \mathbf{\mathcal{Y}}_{C} &\stackrel{\Delta}{=} \mathbf{\mathcal{M}}^{-1} - \mathbf{\mathcal{M}}^{-1} \mathbf{G}_{c}^{*} \left[ \mathbf{G}_{c} \mathbf{\mathcal{M}}^{-1} \mathbf{G}_{c}^{*} \right]^{-1} \mathbf{G}_{c} \mathbf{\mathcal{M}}^{-1} \quad \in \mathbf{\mathcal{R}}^{\mathbf{\mathcal{N}} \times \mathbf{\mathcal{N}}} \end{split}$$

Mass matrix singularity is a consequence of the non-minimal coordinates, DAE approach.





## Augmented Dynamics with Loop Constraints



## Single loop constraint



For loop constraints, the constraint is on the relative motion across bodies

$$Q_x^{rel} \mathcal{V}_x = \mathbf{0}$$
 or  $Q^{rel} [\mathcal{V}_x - \mathcal{V}_y] = \mathbf{0}$ 

With

$$\begin{split} & \mathcal{V}_{nd} \triangleq \begin{bmatrix} \mathcal{V}_{x} \\ \mathcal{V}_{y} \end{bmatrix} & \mathcal{Q} \triangleq \begin{bmatrix} \mathcal{Q}_{x}^{rel}, -\mathcal{Q}_{y}^{rel} \end{bmatrix} \\ & \text{have} & \mathcal{Q}\mathcal{V}_{nd} = \mathbf{0} \end{split}$$



## Loop-constraints can change over time





## **Recall:** B **Pick-Off Operator**



There are times when we need to narrow attention to the nodes



mapping from body to node spatial velocities



## **Recall: Jacobian Matrix**



Combining

$$\mathcal{V}_{nd} \stackrel{3.47}{=} \mathcal{B}^* \mathcal{V}$$
 and  $\mathcal{V} = \Phi^* \mathcal{H}^* \dot{\theta}$ 

#### we have



The <u>Jacobian</u> relates the generalized velocities to the spatial velocity of one or more nodes of interest



## **Constraint matrix with loop constraints**



$$G_{c}(\theta, t)\mathbf{\dot{\theta}} - \mathfrak{U}(t) = \mathbf{0}$$

$$\mathcal{V}_{nd} = \mathcal{J} \dot{\theta}$$
 where  $\mathcal{J} \stackrel{\triangle}{=} \mathcal{B}^* \phi^* \mathcal{H}^*$ 

and 
$$QV_{nd} = 0$$

Hence 
$$G_c = Q\mathcal{J} = Q\mathcal{B}^* \varphi^* H^*$$



### **Recall:** dynamics equations



$$\begin{split} \mathbf{\ddot{\theta}}_{f} &\triangleq \mathcal{M}^{-1} \left( \mathcal{T} - \mathcal{C} \right) \\ \lambda &= - \left[ \mathbf{G}_{c} \mathcal{M}^{-1} \mathbf{G}_{c}^{*} \right]^{-1} \mathbf{\ddot{\delta}}_{f} \\ \mathbf{\ddot{\theta}}_{\delta} &\triangleq \mathcal{M}^{-1} \mathbf{G}_{c}^{*} \lambda \end{split}$$

*Complex, so we look for more simplification* 

$$\ddot{\boldsymbol{\theta}} = \ddot{\boldsymbol{\theta}}_{\mathrm{f}} + \ddot{\boldsymbol{\theta}}_{\delta}$$



## $G_{c}\mathcal{M}^{-1}G_{c}^{\ast}$ operator simplification



Have

Thus,

$$\begin{aligned} \mathbf{G}_{\mathbf{c}} &= \mathcal{Q}\mathcal{J} = \mathcal{Q}\mathcal{B}^{*}\boldsymbol{\varphi}^{*}\mathbf{H}^{*} \qquad [\mathbf{I} - \mathbf{H}\boldsymbol{\psi}\mathcal{K}]\mathbf{H}\boldsymbol{\varphi} = \mathbf{H}\boldsymbol{\psi} \\ \mathbf{G}_{\mathbf{c}}\mathcal{M}^{-1}\mathbf{G}_{\mathbf{c}}^{*} \overset{11.18,9.52}{=} \mathcal{Q}\mathcal{B}^{*}\boldsymbol{\varphi}^{*}\mathbf{H}^{*}[\mathbf{I} - \mathbf{H}\boldsymbol{\psi}\mathcal{K}]^{*}\mathcal{D}^{-1}[\mathbf{I} - \mathbf{H}\boldsymbol{\psi}\mathcal{K}]\mathbf{H}\boldsymbol{\varphi}\mathcal{B}\mathcal{Q}^{*} \\ \overset{9.46}{=} \mathcal{Q}\mathcal{B}^{*}\boldsymbol{\psi}^{*}\mathbf{H}^{*}\mathcal{D}^{-1}\mathbf{H}\boldsymbol{\psi}\mathcal{B}\mathcal{Q}^{*} = \mathcal{Q}\mathcal{B}^{*}\mathcal{\Omega}\mathcal{B}\mathcal{Q}^{*} \overset{10.12}{=} \underline{\mathcal{Q}}\underline{\Lambda}\mathcal{Q}^{*} \end{aligned}$$

where

$$\Omega = \psi^* H^* \mathcal{D}^{-1} \psi H \quad \text{and} \quad \underline{\Lambda} = \mathcal{B}^* \Omega \mathcal{B}$$

QAQ is much much simpler than where we started, but we can do better



## **Operator simplification (contd)**



$$\begin{split} \mathcal{M}^{-1} \mathbf{G}_{\mathbf{c}}^{*} \stackrel{11.18,9.52}{=} [\mathbf{I} - \mathbf{H} \boldsymbol{\psi} \mathcal{K}]^{*} \mathcal{D}^{-1} [\mathbf{I} - \mathbf{H} \boldsymbol{\psi} \mathcal{K}] \mathbf{H} \boldsymbol{\phi} \mathcal{B} \mathcal{Q}^{*} \\ \stackrel{9.46}{=} [\mathbf{I} - \mathbf{H} \boldsymbol{\psi} \mathcal{K}]^{*} \mathcal{D}^{-1} \mathbf{H} \boldsymbol{\psi} \mathcal{B} \mathcal{Q}^{*} \\ [\mathbf{I} - \mathbf{H} \boldsymbol{\psi} \mathcal{K}] \mathbf{H} \boldsymbol{\phi} = \mathbf{H} \boldsymbol{\psi} \\ identity \end{split}$$



## **Simplified equations of motion**



$$\begin{split} \mathbf{\tilde{\theta}}_{f} &\triangleq [\mathbf{I} - \mathbf{H}\psi \mathcal{K}]^{*} \mathcal{D}^{-1} \{ \mathcal{T} - \mathbf{H}\psi [\mathcal{K}\mathcal{T} + \mathcal{P}\mathfrak{a} + \mathfrak{b}] \} - \mathcal{K}^{*}\psi^{*}\mathfrak{a} \\ \lambda &= -[\Omega \underline{\Lambda} \Omega^{*}]^{-1} \mathbf{\tilde{\delta}} \\ \mathbf{\tilde{\theta}}_{\delta} &\triangleq [\mathbf{I} - \mathbf{H}\psi \mathcal{K}]^{*} \mathcal{D}^{-1} \mathbf{H}\psi \mathcal{B} \Omega^{*}\lambda \end{split}$$
$$\begin{aligned} \Omega &= \psi^{*} \mathbf{H}^{*} \mathcal{D}^{-1} \mathbf{\Psi} \mathbf{H} \qquad \underline{\Lambda} = \mathcal{B}^{*} \Omega \mathcal{B} \end{split}$$

Simpler, but still complex, so we continue look for more simplification





## **Simplifying** $\Omega = \psi^* H^* \mathcal{D}^{-1} \psi H$



### **Operational space terminology**



$$\underline{\Lambda} \stackrel{\Delta}{=} \mathcal{J}\mathcal{M}^{-1}\mathcal{J}^{*} \in \mathcal{R}^{6n_{nd} \times 6n_{nd}}$$

$$\underline{\Lambda} \stackrel{\Delta}{=} \underline{\Lambda}^{-1} = (\mathcal{J}\mathcal{M}^{-1}\mathcal{J}^{*})^{-1}$$

$$\overset{\text{aka Operational Space Inertia}}{\text{Matrix (OSIM) in robotics}} \overset{\text{aka Operational Space}}{\text{Compliance Matrix (OSCM)}}$$

$$\overset{\text{aka Extended Operational Space Compliance Matrix (SCM)}}{\overset{\text{simpler operator}}{\text{expression}}} \underbrace{\Omega = \psi^{*} H^{*} \mathcal{D}^{-1} \psi H}$$



## **Why Operational Space? Robotics motivation**

- The operational space (Khatib) is defined as the task space, that is the world viewed from the endeffector that actually interacts with the world
- Generalization to multiple "endeffectors" for legged robots – where legs meet the ground
- Related to area of "whole-body motion control"
- OSCM always exists, but may be singular. Hence the OSIM may not always exist



$$\Lambda = (\mathcal{J}\mathcal{M}^{-1}\mathcal{J}^*)^{-1}$$

Effective task space mass matrix reflected to the nodes of interest







The EOSCM is a mapping from external spatial forces at the nodes to the induced spatial accels at the nodes

$$\Omega \stackrel{\triangle}{=} \psi^* H^* \mathcal{D}^{-1} H \psi \in \mathcal{R}^{6n \times 6n}$$

The EOSCM contains the Backwards Lyapunov form:

$$\mathsf{Z}=\mathbb{A}^*X\mathbb{B}$$





## Generalized Backward Lyapunov Equation for SKO Models





With A & B being SPO operators, and X block diagonal, then







- This product is the dual to the  $Z \triangleq AXB^*$ product used for understanding the mass matrix structure
- Why is this dual product important?
  - It shows up in products of the  $G_c \mathcal{M}^{-1} G_c^*$  form in dynamics analysis
  - One example in cut-joint closed-chain dynamics computations
  - Another example is that of operational space dynamics in robotics



Generalized <u>Backward</u> Lyapunov docomposition

Dual to the forward Lyapunov decomposition

$$\begin{split} \overline{\mathsf{Z}} &= \mathbb{A}^* X \mathbb{B} \qquad \text{block} \\ \text{diagonal} \\ \overline{\mathsf{X}} &= \mathsf{Y} - \operatorname{diagOf}\left\{\mathcal{E}^*_{\mathbb{A}} \mathsf{Y} \mathcal{E}_{\mathbb{B}}\right\} \qquad \overline{\mathsf{Z}} &= \check{\mathsf{Y}} + \tilde{\mathbb{A}}^* \mathsf{Y} + \mathsf{Y} \tilde{\mathbb{B}} + \mathsf{R} \\ \hline \mathsf{R}(\mathfrak{i}, \mathfrak{j}) &= \mathbbm{1}_{[\mathfrak{i} \not\prec \not\succ \mathfrak{j}, \mathsf{k} = \wp(\mathfrak{i}, \mathfrak{j})]} \mathbbm{A}^*(\mathsf{k}, \mathfrak{i}) \ \mathsf{Y}(\mathsf{k}) \ \mathbbm{B}(\mathsf{k}, \mathfrak{j}) \end{split}$$



**Diagonal terms** 



The block-diagonal Y terms can be computed via a O(N) scatter algorithm

$$\begin{split} X &= Y - \operatorname{diagOf}\left\{\mathcal{E}_{\mathbb{A}}^{*}Y\mathcal{E}_{\mathbb{B}}\right\} \\ & [I] \\ Y(k) &= \mathbb{A}^{*}(\wp(k),k)Y(\wp(k))\mathbb{B}(\wp(k),k) + X(k) \end{split}$$

 $\begin{cases} \textbf{for all nodes } k & (base-to-tips \ scatter) \\ Y(k) = \mathbb{A}^*(\wp(k), k) Y(\wp(k)) \mathbb{B}(\wp(k), k) + X(k) \\ \textbf{end loop} \end{cases}$ 



## **General expression for the elements**







### **Derivation**



$$Z = \mathbb{A}^* X \mathbb{B}$$

$$Z(i,j) = \mathbb{A}^*(p,i) \ X(p,q) \ \mathbb{B}(q,j) = \mathbb{A}^*(p,i) \ X(p) \ \mathbb{B}(p,j)$$

$$= \mathbb{1}_{[i \leq p]} \ \mathbb{1}_{[j \leq p]} \ \mathbb{A}^*(p,i) \ X(p) \ \mathbb{B}(p,j)$$

$$I(i,j) = \mathbb{A}^*(k,i) \ Y(k) \ \mathbb{B}(k,j)$$



## **Algorithm Structure**







**Serial-chain case simplification** 



$$\mathsf{Z}=\mathbb{A}^*X\mathbb{B}$$

## All bodies are related in a serial-chain, and hence

$$Z = Y + \tilde{\mathbb{A}}^* Y + Y \tilde{\mathbb{B}} + \mathbb{X}$$
  
$$X = Y - \mathcal{E}^*_{\mathbb{A}} Y \mathcal{E}_{\mathbb{B}}$$
 all zero  
block-diagonal





## Back to Extended Operational Space Compliance Matrix



Applying the Backward Lyapunov decomposition to the EOSCM



$$\begin{split} \Omega &\stackrel{\triangle}{=} \psi^* H^* \mathcal{D}^{-1} H \psi \\ & \checkmark \\ \Omega &= \Upsilon + \tilde{\psi}^* \Upsilon + \Upsilon \tilde{\psi} + R \\ H^* \mathcal{D}^{-1} H &= \Upsilon - \text{diagOf} \left\{ \mathcal{E}^*_{\psi} \Upsilon \mathcal{E}_{\psi} \right\} \end{split}$$

 $\Upsilon(\mathbf{k}) = \psi^*(\wp(\mathbf{k}), \mathbf{k})\Upsilon(\wp(\mathbf{k}))\psi(\wp(\mathbf{k}), \mathbf{k}) + \mathsf{H}^*(\mathbf{k})\mathcal{D}^{-1}(\mathbf{k})\mathsf{H}(\mathbf{k})$ 



## **Component values**



 $\begin{cases} \mbox{for all nodes } k & (\textit{base-to-tips scatter}) \\ \Upsilon(k) = \psi^*(\wp(k), k) \Upsilon(\wp(k)) \psi(\wp(k), k) + H^*(k) \mathcal{D}^{-1}(k) H(k) \\ \mbox{end loop} \end{cases}$ 

$$\Omega(\mathbf{i},\mathbf{j}) = \begin{cases} \Upsilon(\mathbf{i}) & \text{for } \mathbf{i} = \mathbf{j} \\ \Omega(\mathbf{i},\mathbf{k})\psi(\mathbf{k},\mathbf{j}) & \text{for } \mathbf{i} \succeq \mathbf{k} \succ \mathbf{j}, \quad \mathbf{k} = \wp(\mathbf{j}) \\ \Omega^*(\mathbf{j},\mathbf{i}) & \text{for } \mathbf{i} \prec \mathbf{j} \\ \Omega(\mathbf{i},\mathbf{k})\psi(\mathbf{k},\mathbf{j}) & \text{for } \mathbf{i} \neq \mathbf{j}, \quad \mathbf{j} \neq \mathbf{i}, \quad \mathbf{k} = \wp(\mathbf{i},\mathbf{j}) \end{cases}$$



## **Scatter Algorithm Structure**





Another example of low-cost algorithms developed from the SOA methodology. This is the fastest available algorithm to date for the EOSCM.





## For a serial-chain system, the R term is zero. Thus

## $\Omega = \Upsilon + \tilde{\psi}^* \Upsilon + \Upsilon \tilde{\psi}$

$$\mathsf{H}^* \mathcal{D}^{-1} \mathsf{H} = \Upsilon - \mathcal{E}_{\psi}^* \Upsilon \mathcal{E}_{\psi}$$











## Back to Augmented Dynamics with Loop Constraints



Recap







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### Low-cost recursive algorithms



$$\begin{split} \boldsymbol{\tilde{\theta}}_{f} &\triangleq [\mathbf{I} - H\boldsymbol{\psi}\mathcal{K}]^{*}\mathcal{D}^{-1}\big\{\mathcal{T} - H\boldsymbol{\psi}[\mathcal{K}\mathcal{T} + \mathcal{P}\boldsymbol{\mathfrak{a}} + \boldsymbol{\mathfrak{b}}]\big\} - \mathcal{K}^{*}\boldsymbol{\psi}^{*}\boldsymbol{\mathfrak{a}} \\ \lambda &= -[\Omega\underline{\Lambda}\Omega^{*}]^{-1}\boldsymbol{\mathfrak{F}} \longleftarrow \overset{\text{EOSCM recursive}}{\underset{computation}{}} \overset{O(N) \text{ ATBI forward}}{\underset{dynamics}{}} \\ \boldsymbol{\tilde{\theta}}_{\delta} &\triangleq [\mathbf{I} - H\boldsymbol{\psi}\mathcal{K}]^{*}\mathcal{D}^{-1}H\boldsymbol{\psi}\mathcal{B}\Omega^{*}\boldsymbol{\lambda} & & \\ \Omega &= \boldsymbol{\psi}^{*}H^{*}\mathcal{D}^{-1}\boldsymbol{\psi}H & \underline{\Lambda} &= \mathcal{B}^{*}\Omega\mathcal{B} \end{split}$$



## **Augmented method comments**



- Even though the augmented method approach does not lend itself directly to be SKO model, we find that the SKO algorithms can be used to efficiently carry out each of the augmented method steps
- The <u>minimal</u> cuts augmented method is much better than the <u>maximal</u> cuts approach
  - We can take advantage of the fast SKO method
  - Smaller number of constraints to manage error for
- Changing of constraints are easily accommodated by the SKO gather & scatter algorithms
- The augmented approach still remains a <u>non-minimal</u> coordinates and <u>DAE</u> approach
  - Hence some type of error control (eg. Baumgarte, projection, implicit method) is needed when integrating





## Recap







- Looked into the augmented method for closed-chain dynamics (DAE approach)
- Does not have a direct SKO model
- Introduced the notion of operational space inertia matrix (OSIM) and OSCIM
- Discussed the Backward Lyapunov Equation based operator decomposition
- Applied SKO model recursive algorithms for the various steps in the augmented approach



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